Superconductivity beyond the dimer model in 2D organic charge transfer salts

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Organic charge transfer salts: Crystal structure and properties of $(ET)_2X$

- ET = BEDT-TTF
 - = bis(ethylenedithio)tetrathiafulvalene is the electron donor
- X is the electron acceptor (e.g. Cu(NCS)₂)
- we concentrate on κ-(ET)₂X salts
- AFI to SC transition with pressure or variation of X



Organic charge transfer salts: Electronic structure of κ -(ET)₂X

- two donor molecules transfer in total one electron to the acceptor X
- two-dimensional electronic structure in the ET-plane
- four ET molecules in the unit cell, 3/4 filled four-band model
- effective model only describes dimers of ET molecules, 1/2 filled two-band model



Figure : Ferber, Foyevtsova, Jeschke, Valentí, PRB **89**, 205106 (2014)



Symmetry of superconducting pairing in κ -(ET)₂X

Ab-initio calculations

 full potential local orbital (FPLO) code
molecular orbital models from projective Wannier functions

Superconductivity calculations

RPA spin-fluctuation pairing

Bickers, Scalapino, White, PRL 62, 961 (1989)

Graser, Maier, Hirschfeld, Scalapino, New J. Phys. 11, 025016 (2009)

 extract pairing symmetry and relative strength





 d_{xy} from experiment and theory 4 /

Competing pairing symmetries in κ -(ET)₂X



see also Schmalian, PRL 81, 4232 (1998); Kuroki et al., PRB 65, 100516(R) (2002)

Material dependence of T_c in κ -(ET)₂X

- λ measures pairing strength
- U is material dependent
- feature rich spin-susceptibility not discussed in previous work





Summary

- we derived ab-initio models for many κ-phase materials
- RPA spin-fluctuation pairing yields $s_{\pm}/d_{x^2-y^2}$ or d_{xy} gap
- material dependence of T_c comes out correctly
- no simple dependence on microscopic parameters

Altmeyer, Guterding, Jeschke, Valentí, to be published

Conversion from four-band to dimer model





Tight binding+RPA formalism in a nutshell

$$\begin{split} \chi_{st}^{pq}(\vec{q}) &= -\frac{1}{N} \sum_{\vec{k},\mu,\nu} a_{\mu}^{s}(\vec{k}) a_{\mu}^{p*}(\vec{k}) a_{\nu}^{q}(\vec{k}+\vec{q}) a_{\nu}^{t*}(\vec{k}+\vec{q}) \frac{f(E_{\nu}(\vec{k}+\vec{q})) - f(E_{\mu}(\vec{k}))}{E_{\nu}(\vec{k}+\vec{q}) - E_{\mu}(\vec{k})} \\ & \left[(\chi_{spin}^{RPA})_{st}^{pq} \right]^{-1} = [\chi_{st}^{pq}]^{-1} - (U_{spin})_{st}^{pq} \\ \Gamma_{st}^{pq}(\vec{k},\vec{k}') &= \left[\frac{3}{2} U_{s} \, \chi_{s}^{RPA}(\vec{k}-\vec{k}') \, U_{s} + \frac{1}{2} U_{s} - \frac{1}{2} U_{c} \, \chi_{c}^{RPA}(\vec{k}-\vec{k}') U_{c} + \frac{1}{2} U_{c} \right]_{ps}^{tq} \\ \Gamma_{ij}(\vec{k},\vec{k}') &= \sum_{stpq} a_{i}^{t*}(-\vec{k}) a_{i}^{s*}(\vec{k}) \text{Re} \left[\Gamma_{st}^{pq}(\vec{k},\vec{k}') \right] a_{j}^{p}(\vec{k}') a_{j}^{q}(-\vec{k}') \\ -\sum_{j} \oint_{C_{j}} \frac{dk'_{\parallel}}{2\pi} \frac{1}{4\pi\nu_{F}(\vec{k}')} \left[\Gamma_{ij}(\vec{k},\vec{k}') + \Gamma_{ij}(\vec{k},-\vec{k}') \right] g_{j}(\vec{k}') = \lambda_{i}g_{i}(\vec{k}) \end{split}$$

Graser, Maier, Hirschfeld, Scalapino, New Journal of Physics 11, 025016 (2009)

Other interesting talks

- talk after next one : Electronic structure of a dual-layered organic charge transfer salt, Harald O. Jeschke
- Z5.00004 : Generalized bandstructure unfolding method, Friday 11:15 AM, Milan Tomić

Summary of superconductivity calculations

- feature rich spin-susceptibility
- RPA spin-fluctuation pairing yields $s_{\pm}/d_{x^2-y^2}$ or d_{xy} gap
- material dependence of T_c comes out correctly
- complicated dependence on microscopic parameters

